# Inverse Problems of Nonlinear Wave Equations 

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- Let $(M, g)$ be a Lorentzian $(n+1)$-manifold of signature $(-,+\cdots,+)$. Consider

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\begin{equation*}
\square_{g} u+N(u)=f,\left.\quad u\right|_{t \leq 0}=0, \tag{1}
\end{equation*}
$$

where $N(\cdot)$ is nonlinear and $\square_{g}:=d^{*} d$ locally reads

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- Let $G$ be a compact Lie group with Lie algebra $\mathfrak{g}$. Given a connection 1-form $A \in C^{\infty}\left(M ; T^{*} M \otimes \mathfrak{g}\right)$, consider

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\begin{equation*}
\square_{g, A} u+N(u)=f,\left.\quad u\right|_{t \leq 0}=0 \tag{2}
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$$

where $\square_{g, A}:=(d+A)^{*}(d+A)$ on $\mathbb{M}^{3+1}$ locally reads

$$
\square_{g, A} u=\square_{g} u-2\left(-A_{0} \frac{\partial u}{\partial t}+\sum_{j=1}^{3} A_{j} \frac{\partial u}{\partial x^{j}}\right)+\left(\operatorname{div} A+A_{0}^{2}-\sum_{j=1}^{3} A_{j}^{2}\right) u
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- Given a time-dependent potential $V$, consider

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- Bony, Reed-Rauch, Melrose-Ritter : nonlinear waves are approximately the interactions of linear waves.
- Melrose-Uhlmann, Guillemin-Uhlmann, Greenleaf-Uhlmann : Such interactions are Fourier Integral Operators associated multiple intersecting Lagrangians.

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- Feizmohammadi-Oksanen : recover $V$ in (3)
- Feizmohammadi-Lassas-Oksanen : introduce the 3-to-1 scattering relation based on the third order linearization to recover $g$ in (1).


## Third order linearisation

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- To linearise the semilinear wave equation,

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we choose

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with sources $f_{(j)}$ around $x_{(j)}$ and small parameters $\epsilon_{(j)}>0$.

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- The derivatives

$$
\begin{aligned}
v_{(j)} & :=\left.\partial_{\epsilon_{(j)}} \phi\right|_{\epsilon=0} \\
v_{(j k)} & :=\left.\partial_{\epsilon_{(j k)}} \phi\right|_{\epsilon=0} \\
v_{(123)} & :=\left.\partial_{\epsilon_{(1)}} \partial_{\epsilon_{(2)}} \partial_{\epsilon_{(3)}} \phi\right|_{\epsilon=0}
\end{aligned}
$$

satisfy linear wave equations

$$
\begin{align*}
\left(\square_{g}-A \cdot d\right) v_{(j)} & =f_{(j)}  \tag{4}\\
\left(\square_{g}-A \cdot d\right) v_{(j k)} & =0  \tag{5}\\
\left(\square_{g}-A \cdot d\right) v_{(123)} & =\frac{1}{2} \sum_{\{i, j, k\}=\{1,2,3\}}\left\langle v_{(i)}, v_{(j)}\right\rangle v_{(k)} \tag{6}
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- From the measurable data $L^{g, A}$, we reconstruct the geometry $(g, A)$ by understanding the X-ray transform $\mathbf{S}^{g, A}$.
Measuable data $\Rightarrow$ X-ray transform $\Rightarrow$ Geometry $L^{g_{1}, A_{1}}=L^{g_{2}, A_{2}} \quad \Rightarrow \quad \mathbf{S}^{g_{1}, A_{1}}=\mathbf{S}^{g_{2}, A_{2}} \quad \Rightarrow \quad\left[g_{1}, A_{1}\right]=\left[g_{2}, A_{2}\right]$,


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Theorem
For any $0<\theta \leq 1$ and $\bar{s}>3$ with $\bar{s} / \theta \in \mathbb{Z}$,

$$
\|[A]-[B]\|_{C_{0}(\mathbb{D} \backslash \mho)} \lesssim C_{A, B} \tilde{C}_{A, B, \theta, \bar{s}}\|A-B\|_{C_{0}^{\bar{s} / \theta}(\mathbb{D})}^{\theta}\left\|L_{A}-L_{B}\right\|_{\mho, s}^{1-\theta} .
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- $C$ : the reconstruction of $A$ in (2) is stable in the sense

Theorem
For any $0<\theta \leq 1$ and $\bar{s}>3$ with $\bar{s} / \theta \in \mathbb{Z}$,

$$
\|[A]-[B]\|_{C_{0}(\mathbb{D} \backslash \mho)} \lesssim C_{A, B} \tilde{C}_{A, B, \theta, \bar{s}}\|A-B\|_{C_{0}^{\bar{s} / \theta}(\mathbb{D})}^{\theta}\left\|L_{A}-L_{B}\right\|_{\mho, s}^{1-\theta} .
$$

- C : Reduction via the third order linearization and the 3-to-1 scattering relation to the stability of the broken X-ray
Theorem

$$
\left\|\mathrm{S}_{z_{\min } \leftarrow \boldsymbol{y} \leftarrow \boldsymbol{x}}^{A}-\mathrm{S}_{z_{\min } \leftarrow \boldsymbol{y} \leftarrow \boldsymbol{x}}^{B}\right\|_{C_{0}\left(\mathcal{F}^{x}\right)} \lesssim\left\|L_{A}-L_{B}\right\|_{\mho, s}
$$

$$
\|\rho([A])-\rho([B])\|_{C_{0}(\mathbb{D} \backslash \mho)} \lesssim \tilde{C}_{A, B}\left\|\mathbf{S}_{z \leftarrow \boldsymbol{y} \leftarrow \boldsymbol{x}}^{A}-\mathbf{S}_{z \leftarrow \boldsymbol{y} \leftarrow \boldsymbol{x}}^{B}\right\|_{C_{0}^{1}\left(\mathcal{F}^{X}\right)}
$$

Thank you for your attention.

