Inverse Problems of Nonlinear Wave Equations

Chen Xi Shanghai Centre for Mathematical Sciences Fudan University

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Let (M,g) be a Lorentzian (n + 1)-manifold of signature (-,+...,+). Consider

$$\Box_{g} u + N(u) = f, \quad u|_{t \le 0} = 0,$$
 (1)

where $\mathit{N}(\cdot)$ is nonlinear and $\Box_{g}:=d^{*}d$ locally reads

$$\Box_g u = (-\det(g))^{-1/2} \partial_{x^{\alpha}} \Big((-\det(g))^{-1/2} g^{\alpha\beta} \partial_{x^{\beta}} u(x) \Big).$$

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Let G be a compact Lie group with Lie algebra g. Given a connection 1-form A ∈ C[∞](M; T*M ⊗ g), consider

$$\Box_{g,A} u + N(u) = f, \quad u|_{t \le 0} = 0,$$
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where $\Box_{g,\mathcal{A}} := (d+\mathcal{A})^*(d+\mathcal{A})$ on \mathbb{M}^{3+1} locally reads

$$\Box_{g,A} u = \Box_g u - 2 \left(-A_0 \frac{\partial u}{\partial t} + \sum_{j=1}^3 A_j \frac{\partial u}{\partial x^j} \right) + \left(\operatorname{div} A + A_0^2 - \sum_{j=1}^3 A_j^2 \right) u.$$

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Given a time-dependent potential V, consider

$$\Box_{g,A}u + Vu + N(u) = f, \quad u|_{t \le 0} = 0. \tag{3}$$

► General Relativity:

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 - The wave equation (1) describes the wave propagation and asymptotic in the spacetime (M, g) (e.g. black holes) and the source f could be gravitational waves.

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▶ The Klein-Gordon equation on \mathbb{M}^{3+1} with a potential :

$$\Box u + m^2 u + V u + N(u) = 0.$$

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The inverse problem : recover (g, A, V) from reasonable local measurements in relevant physical models.

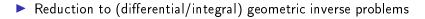
- The inverse problem : recover (g, A, V) from reasonable local measurements in relevant physical models.
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the Source-to-Solution map : $f \mapsto u|_{\mho}$

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► the X-ray inversion ⇒ the potential V

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Microlocal analysis of nonlinear wave equations :

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 - <u>Melrose-Uhlmann</u>, <u>Guillemin-Uhlmann</u>, <u>Greenleaf-Uhlmann</u>: Such interactions are Fourier Integral Operators associated multiple intersecting Lagrangians.

Progresses II

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- Feizmohammadi-Oksanen : recover V in (3)
- Feizmohammadi-Lassas-Oksanen : introduce the 3-to-1 scattering relation based on the third order linearization to recover g in (1).

Third order linearisation

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► To linearise the semilinear wave equation,

$$\Box_{g}\phi - A \cdot d\phi + |\phi|^{2}\phi = f, \quad \phi|_{t<0} = 0,$$

we choose

$$f = \epsilon_{(1)}f_{(1)} + \epsilon_{(2)}f_{(2)} + \epsilon_{(3)}f_{(3)}$$

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with sources $f_{(j)}$ around $x_{(j)}$ and small parameters $\epsilon_{(j)} > 0$.

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$$f = \epsilon_{(1)}f_{(1)} + \epsilon_{(2)}f_{(2)} + \epsilon_{(3)}f_{(3)}$$

with sources *f*_(*j*) around *x*_(*j*) and small parameters *ϵ*_(*j*) > 0. ► The derivatives

$$\begin{aligned} \mathbf{v}_{(j)} &:= \partial_{\epsilon_{(j)}} \phi|_{\epsilon=0} \\ \mathbf{v}_{(jk)} &:= \partial_{\epsilon_{(jk)}} \phi|_{\epsilon=0} \\ \mathbf{v}_{(123)} &:= \partial_{\epsilon_{(1)}} \partial_{\epsilon_{(2)}} \partial_{\epsilon_{(3)}} \phi|_{\epsilon=0} \end{aligned}$$

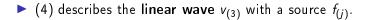
satisfy linear wave equations

$$(\Box_g - A \cdot d)v_{(j)} = f_{(j)}$$
(4)

$$(\Box_g - A \cdot d) v_{(jk)} = 0$$
⁽⁵⁾

$$(\Box_g - A \cdot d) v_{(123)} = \frac{1}{2} \sum_{\{i,j,k\} = \{1,2,3\}} \langle v_{(i)}, v_{(j)} \rangle v_{(k)}.$$
 (6)

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- (4) describes the **linear wave** $v_{(3)}$ with a source $f_{(j)}$.
- The principal symbol satisfies a parallel transport equation

$$\mathcal{L}_{H_p}\sigma(\mathsf{v}_{(j)}) + \langle \mathsf{A}, \dot{\gamma} \rangle \sigma(\mathsf{v}_{(j)}) = \nabla^{\mathsf{A}}_{\dot{\gamma}}\sigma(\mathsf{v}_{(j)}) = 0 \quad \text{along geodesics } \gamma.$$

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▶ RHS of (6) is the interaction of linear waves $v_{(1)}, v_{(2)}, v_{(3)}$.

- (4) describes the linear wave $v_{(3)}$ with a source $f_{(j)}$.
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$$\mathcal{L}_{H_{p}}\sigma(\mathsf{v}_{(j)}) + \langle A, \dot{\gamma} \rangle \sigma(\mathsf{v}_{(j)}) = \nabla^{A}_{\dot{\gamma}}\sigma(\mathsf{v}_{(j)}) = 0 \quad \text{along geodesics } \gamma.$$

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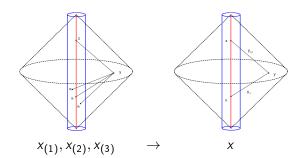
RHS of (6) is the interaction of linear waves v₍₁₎, v₍₂₎, v₍₃₎.
 v₍₁₂₃₎ is a returning linear wave created by the interaction.

Nonlinear interaction of three waves

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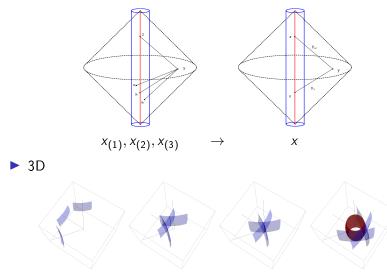
Nonlinear interaction of three waves

▶ 2D



Nonlinear interaction of three waves

▶ 2D



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Consider the perturbed semilinear wave equation

$$\Box_g \phi - A \cdot d\phi + |\phi|^2 \phi = f, \quad \phi|_{t<0} = 0.$$

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From the measurable data $L^{g,A}$, we reconstruct the geometry (g, A) by understanding the X-ray transform $S^{g,A}$.

 $\begin{array}{rcl} \text{Measuable data} & \Rightarrow & \text{X-ray transform} & \Rightarrow & \text{Geometry} \\ L^{g_1,A_1} = L^{g_2,A_2} & \Rightarrow & \textbf{S}^{g_1,A_1} = \textbf{S}^{g_2,A_2} & \Rightarrow & [g_1,A_1] = [g_2,A_2], \\ & & & & & & \\ \end{array}$

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The inverse problem : What happens to the result if there are measurement errors such as noises ?

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C : the reconstruction of A in (2) is stable in the sense

Theorem

For any $0 < \theta \leq 1$ and $\bar{s} > 3$ with $\bar{s}/\theta \in \mathbb{Z}$,

$$\|[A]-[B]\|_{\mathcal{C}_0(\mathbb{D}\setminus\mho)} \lesssim C_{A,B} \tilde{\mathcal{C}}_{A,B,\theta,\overline{s}} \|A-B\|_{\mathcal{C}_0^{\overline{s}/\theta}(\mathbb{D})}^{\theta} \|L_A-L_B\|_{\mho,s}^{1-\theta}.$$

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 C : Reduction via the third order linearization and the 3-to-1 scattering relation to the stability of the broken X-ray

Theorem

$$\left\| \mathbf{S}_{z_{\min} \leftarrow \mathbf{y} \leftarrow \mathbf{x}}^{A} - \mathbf{S}_{z_{\min} \leftarrow \mathbf{y} \leftarrow \mathbf{x}}^{B} \right\|_{C_{0}(\mathcal{F}^{X})} \lesssim \left\| L_{A} - L_{B} \right\|_{\mathcal{U},s},$$
$$\left\| \rho([A]) - \rho([B]) \right\|_{C_{0}(\mathbb{D} \setminus \mathcal{U})} \lesssim \tilde{C}_{A,B} \left\| \mathbf{S}_{z \leftarrow \mathbf{y} \leftarrow \mathbf{x}}^{A} - \mathbf{S}_{z \leftarrow \mathbf{y} \leftarrow \mathbf{x}}^{B} \right\|_{C_{0}^{1}(\mathcal{F}^{X})}.$$

Thank you for your attention.